

# Impact Deflection of Potentially Hazardous Asteroids Using Current Launch Vehicles

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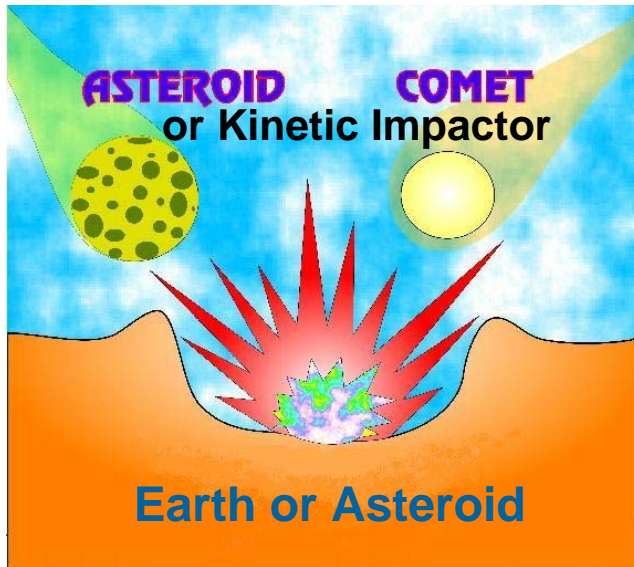
## Why Kinetic Impact?

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- ✓ Simple
- ✓ Inexpensive (relative to other methods)
- ✓ Fastest to effect
- ✓ Requires no new or exotic technology
- ✓ Doesn't require building or testing nuclear warheads
- ✓ Highly capable
- ✓ We can do it NOW.

## Why *Call* it Kinetic Impact?

- Kinetic Impact (like all deflection methods) involves momentum transfer, so why call it Kinetic Impact? Why not Momentum Impact?
  - Just like asteroid striking Earth, impactor striking asteroid has huge kinetic energy which, upon impact, causes explosive cratering.
  - Ejecta from crater acts as propulsion.



## Study Overview

- We model kinetic impact missions to 795 PHAs taken from [neo.jpl.nasa.gov](http://neo.jpl.nasa.gov) on 6/21/06.
- All of our simulations use the actual first four orbital elements [a e i  $\omega$ ] for the all 795 asteroids. We assume a circular Earth orbit, and no knowledge of mission date.
- Other than orbits, we take a parametric approach, varying:
  - Asteroid size
  - Asteroid density
  - Cratering parameters
  - Warning time

## Deflection Distance

- Start with a cube-shaped Earth, and an asteroid heading toward the center of an Earth face at  $90^\circ$  and Earth velocity.
  - Required deflection distance is  $1 R_\oplus$
- Factors that decrease required deflection distance:
  - Earth is sphere-like
  - Angle between Earth and asteroid velocities is  $< 90^\circ$
- Factor that increases required deflection distance:
  - Gravitational focusing makes Earth effectively  $\sim 1.4$  times bigger at typical Earth-asteroid relative velocities

We use  $1 R_\oplus$  in our simulations.

(Distance is  $\sim 4$  orders of magnitude less for keyholes.)

# Orbit Perturbation

- In-track perturbation gives cumulative effect over time.
- For extremely short warning times (< 1 year), optimal perturbation may not be in-track, but for likely warning times, it will be almost purely in-track.

→ We use in-track perturbation only, with two body model:

$$\frac{da}{dt} = \frac{2a^{3/2}}{\sqrt{\mu(1-e^2)}} \left[ \frac{dV}{dt} (1 + e \cdot \cos v) \right]$$

$$dV = \frac{2\pi \cdot \delta \cdot a \cdot (1-e^2)^{0.5}}{3 \cdot T \cdot C \cdot [e \sin \phi \sin v + K_\phi (1 + e \cos v)]}$$

$$\phi = \cos^{-1} K_\phi$$

$$K_\phi = \left[ \frac{(1 + e \cos v)^2}{2(1 + e \cos v) - (1 - e^2)} \right]^{0.5}$$

$$C \approx \pi(a+b) \left[ 1 + \frac{3x^2}{10 + \sqrt{4 - 3x^2}} \right] \quad x = \frac{a-b}{a+b}$$

**a = semi-major axis**

**b = semi-minor axis**

**V = velocity**

**μ = sun grav. constant**

**e = eccentricity**

**v = true anomaly**

**φ = flight path angle**

**δ = deflection distance**

**T = warning time**

**C = orbit circumference**

## Impact Location on Asteroid

- Assume that ejecta average momentum is normal to surface impacted.
- Impactor should strike asteroid at location where surface is normal to asteroid velocity vector, to give in-track  $\Delta V$ .
- Imparting spin to asteroid is ok because linear and angular momentum are conserved separately.
- Final impactor targeting will not be trivial, but is doable with current technology.
  - See Deep Impact, MDA kinetic kill vehicles

## Crater Dynamics

- Magnitude of ejecta momentum depends on cratering dynamics.
- We draw on work of Holsapple, Housen, Schmidt, and Melosh to model cratering event.
- Crater model is influenced strongly by three target dependant parameters:  $K$ ,  $C_D$ ,  $\zeta$  ( $\zeta$  is often called  $\nu$  in literature)
- We use a parameter set that was determined experimentally for dry Ottawa sand:

$$K = 0.32 \quad C_D = 1.68 \quad \zeta = 1.22$$

- This is a conservative set, i.e. gives low momentum ratios compared to others in literature.
- Also gives low momentum ratios compared to other cratering models.

# Ejecta Momentum

$$\Delta P_{ej} = \int_{V_{esc}}^{\infty} V_{final} \cdot \sin \theta \cdot \frac{dm_{ej}}{dV_{ej}} \cdot dV_{ej}$$

$$V_{final} = \sqrt{V_{ej}^2 - V_{esc}^2}$$

$\frac{dm_{ej}}{dV_{ej}}$  is derived from:

$$Vol_{ej} (> V_{ej}) R^{-3} = K \left( \sqrt{gR} / V_{ej} \right)^{\zeta}$$

(Use gravity scaling.)

R is derived from other crater relations in the literature.

$\Delta P_{ej}$  = momentum imparted by ejecta

$v_{esc}$  = asteroid escape velocity

$v_{final}$  = ejecta vel. after escape

$\theta$  = ejecta angle (assumed =  $45^\circ$ )

$V_{ej}$  = ejecta velocity

$m_{ej}$  = ejecta mass

$V_i$  = impact velocity

$m_{ast}$  = asteroid mass

$Vol_{ej}$  = ejecta velocity

$\zeta$  = cratering constant

$K$  = cratering constant

# Impactor Mass Calculation

$$m_i = \left( -B - \sqrt{B^2 - 4AC} \right) / 2A$$

where

$$A = k_1 k_4$$

$$B = -\Delta V \cdot k_3 k_4 - k_1 k_3 - k_2 k_3$$

$$C = \Delta V \cdot k_3^2$$

and

$$k_1 = V_i |\sin \omega_i|$$

$$k_2 = \sin \theta \cdot \zeta K C_D^{3+\zeta/2} 2^{\frac{3(1+\beta)}{\beta-1}} V_i^\zeta \rho_a^{-\zeta/6} \rho_i^{\zeta/6} \cdot \int_{V_{esc}}^{\infty} \left( V_{ej}^2 - 2\mu_a/r_a \right)^{0.5} V_{ej}^{-\zeta-1} dV_{ej}$$

$$k_3 = m_{ast}$$

$$k_4 = \frac{1}{2} K C_D^{3+\zeta/2} 2^{\frac{3(1+\beta)}{\beta-1}} V_i^\zeta \rho_a^{-\zeta/6} \rho_i^{\zeta/6} V_{esc}^{-\zeta}$$

$V_i$  = impact velocity

$\omega_i$  = impact angle

$\rho_i$  = impactor density

$\rho_a$  = asteroid density

$m_{ast}$  = asteroid mass

$\theta$  = ejecta angle (assumed = 45°)

$V_{esc}$  = asteroid escape velocity

$V_{ej}$  = ejecta velocity

$\zeta$  = cratering constant

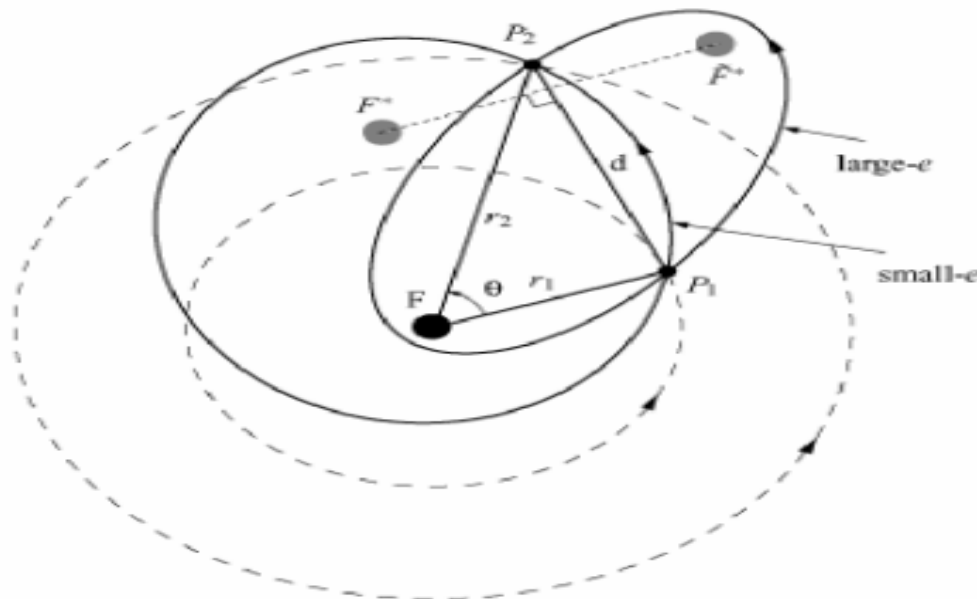
$K$  = cratering constant

$C_D$  = cratering constant

$\beta$  =  $\zeta / (\zeta + 6)$

# Trajectories

- In real mission, some JPL-types would find optimal interplanetary trajectory, maybe using gravity assist fly-by's.
- To model large number of missions, and to be conservative, we use Lambert trajectories.
- For each given departure and arrival point, simulations analyze all four trajectories for a range of intercept semi-major axes.



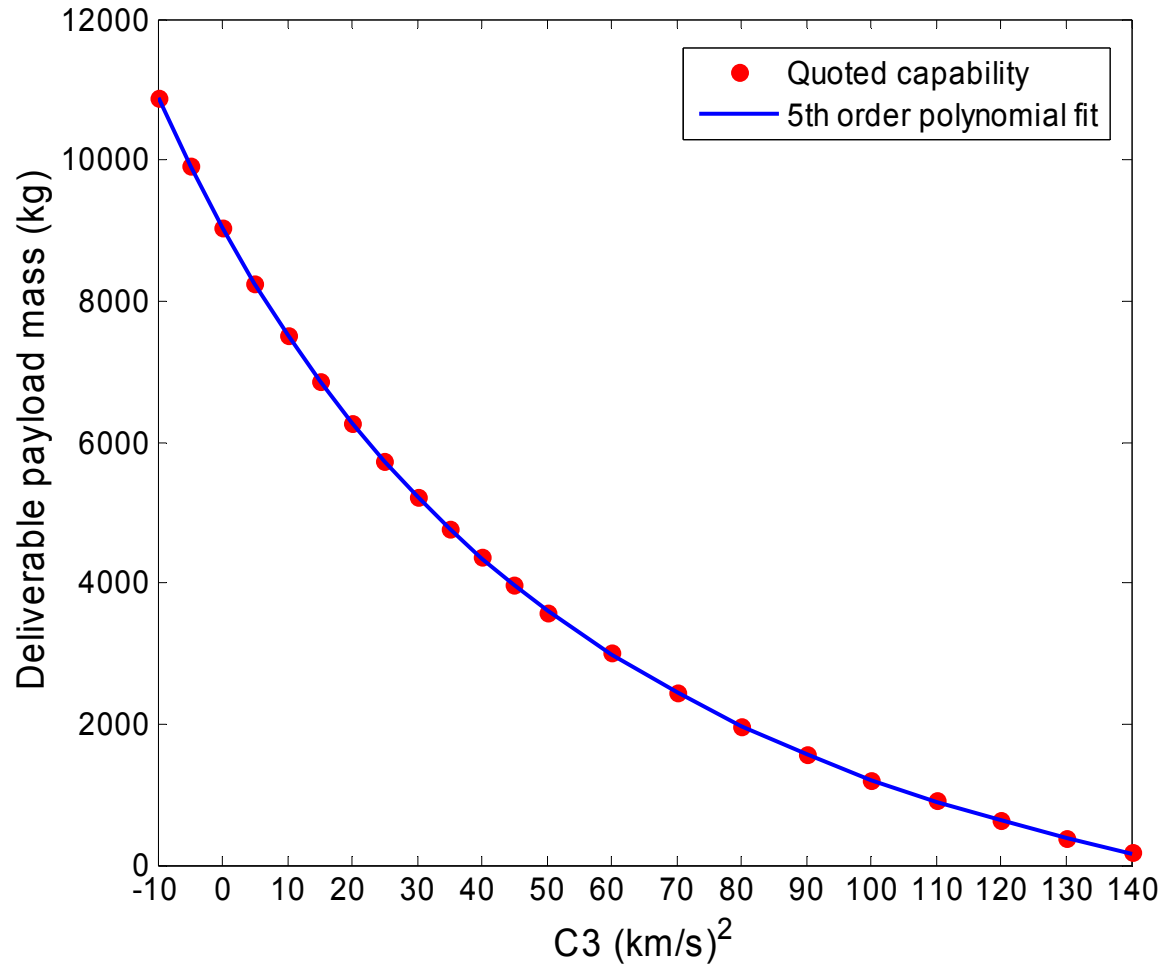
## Earth-Asteroid Geometry

- We could calculate the most efficient deflection theoretically possible by allowing the impactor to depart any point in Earth orbit and arrive at any point in asteroid orbit.
- However, in reality, the resulting optimal geometry might not be present during the time frame in which we want to launch our mission.
- Instead, we use a separate Monte-Carlo analysis to select Earth-asteroid geometries for which we have a certain probability of achieving that geometry or better within a certain time frame.
- There is a 60% chance of achieving our results or better within a 10 year launch window, 80% within 20 years.

*(Remember, this assumes we are limited to simple Lambert trajectories.)*

# Launch Vehicle

- We assume only present launch vehicle capability. (Not including Ares V)
- Our model uses Atlas V Heavy Lift Version, Single Engine Centaur upper stage.
- Note: There is an option to leave SEC attached to payload. This adds another ~2500 kg to delivered impactor mass. We do not include this in our model, but instead leave it as a source of margin for our results.



## Parameterization of Results

- For every intercept trajectory analyzed:
  - Compute required impactor mass for specified deflection.
  - Compute max possible payload mass (using fit function of quoted launch vehicle performance) for the associated C3.

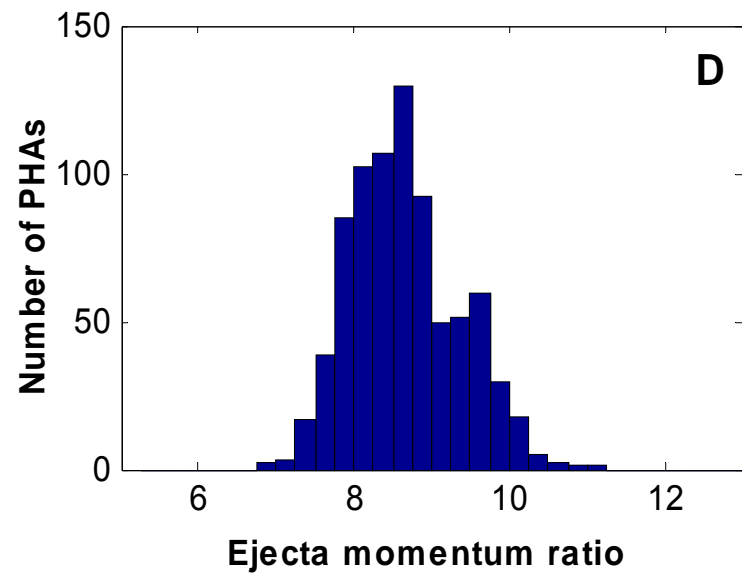
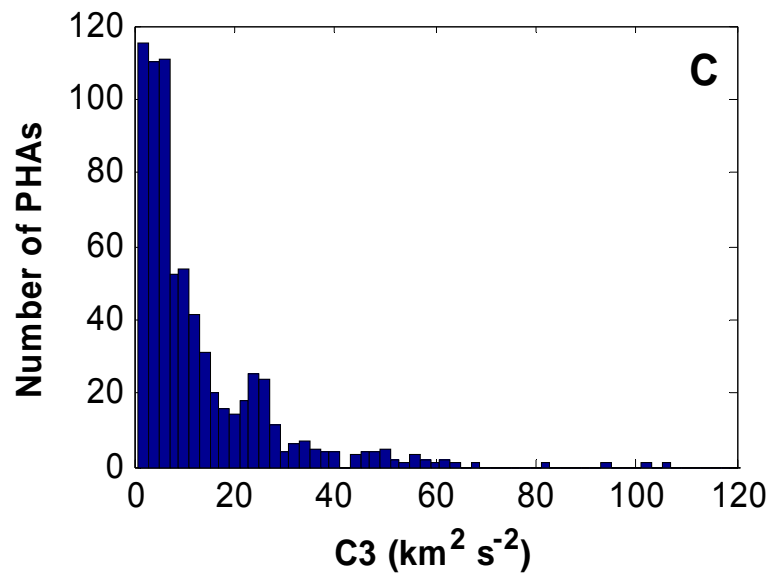
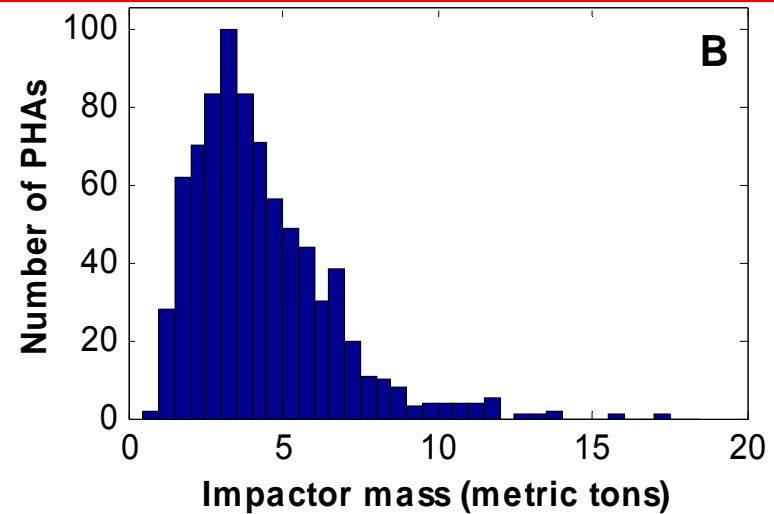
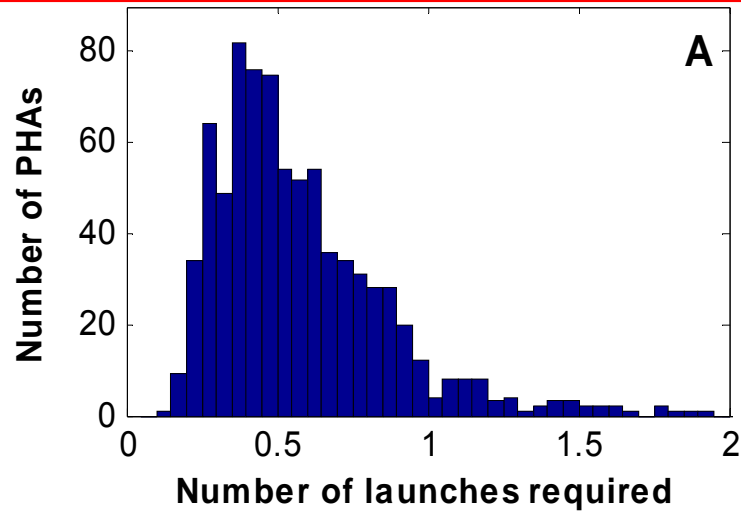
$$\frac{\text{Required Mass}}{\text{Deliverable Mass}} = \lambda_{mass}$$

- $\lambda_{mass}$  is our metric to optimize, and to evaluate the viability of kinetic impact method.
- Rounding up  $\lambda_{mass}$  to the next integer gives the number of launches required.

# Results

Simulation	Asteroid diameter (m)	Asteroid density (g cm <sup>-3</sup> )	Impactor density (g cm <sup>-3</sup> )	Time from deflection impact to predicted Earth collision (years)	Crater model	Number of launches / Percentage of PHAs that can be deflected with that number of launches or fewer	
1	1000	3	19	100	Ottawa sand	1 / 55%	2 / 94%
2	1000	3	19	20	Ottawa sand	5 / 55%	10 / 94%
3	500	3	19	20	Ottawa sand	1 / 93%	2 / 100%
4	250	3	19	5	Ottawa sand	1 / 100%	
5	1000	3	8.9	100	Ottawa sand	1 / 44%	2 / 91%
6	1000	1.5	19	100	Ottawa sand	1 / 97%	2 / 100%
7	1000	0.5	19	20	Ottawa sand	1 / 95%	2 / 100%
8	500	1.5	19	10	Ottawa sand	1 / 100%	
9	500	0.5	19	3	Ottawa sand	1 / 99%	2 / 100%
10	1000	3	NA	50	P ratio = 38.5	1 / 99%	2 / 100%
11	1000	3	NA	100	No ejecta	10 / 57%	30 / 97%
12	750	0.5	NA	100	No ejecta	1 / 74%	2 / 95%
13	400	3	NA	100	No ejecta	1 / 78%	2 / 97%
14	500	0.5	NA	20	No ejecta	1 / 54%	3 / 96%
15	250	3	NA	20	No ejecta	1 / 71%	3 / 99%
16	250	1.5	NA	20	No ejecta	1 / 92%	2 / 100%
17	250	0.5	NA	5	No ejecta	1 / 86%	2 / 99%
18	150	3	NA	20	No ejecta	1 / 100%	

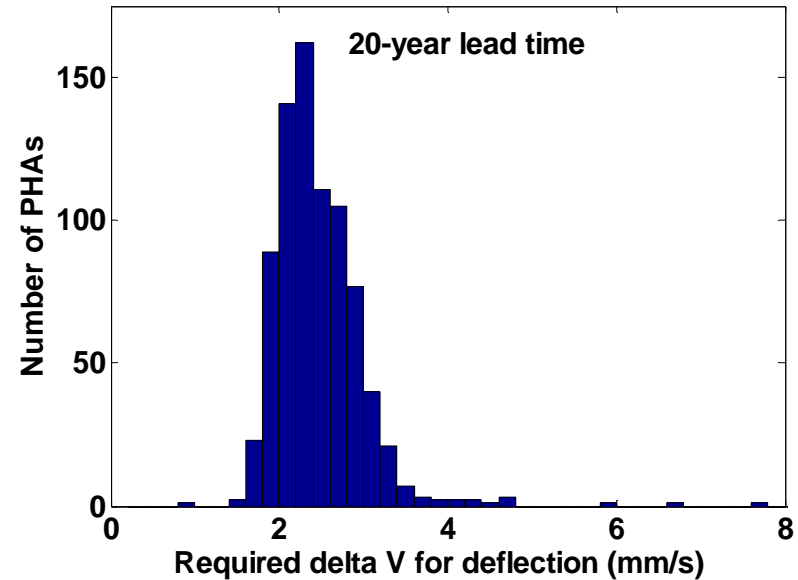
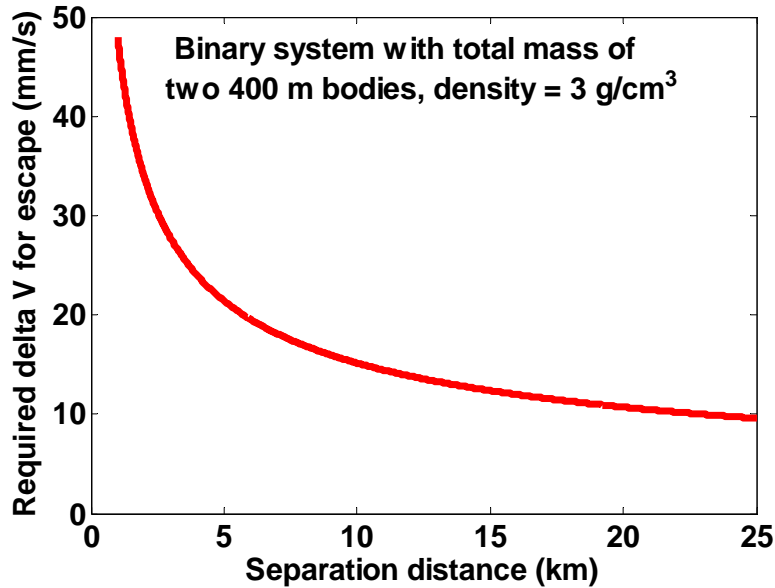
# Results for Sim 3 (500 m, 3 g/cm<sup>3</sup>, 20 years)



## Binary Asteroids

- 3 cases:
  - A)  $\Delta V$  is large enough and in precise direction to permanently separate two bodies
  - B) Binary orbit is perturbed such that bodies stick together, forming “contact binary”
  - C) The two bodies remain in orbit around each other
- In Cases B and C, momentum is conserved, and the impact has the same effect as on a single body of similar mass.
- Case A requires  $\Delta V$  at least as big as the difference between the bodies' relative velocity,  $V_{rel} = \sqrt{G(m_1 + m_2)/\text{distance}}$ , and their mutual escape velocity,  $V_{eM} = \sqrt{2} \cdot V_{rel}$ .
  - For a given total mass, Case A is less likely for greater size disparity between the two bodies, assuming an impact to the larger body.

# Binary Asteroids



Consider binary with two 400 m bodies (~ mass of single 500 m body):

- At 23 km separation, minimum  $\Delta V$  for mutual escape is 10 mm/s.
- This  $\Delta V$  must be parallel to impacted body's velocity relative to CoM.
- 10 mm/s for one body corresponds to 5 mm/s for the system.
- In vast majority of deflection cases (for lead time  $\geq 20$  years), required  $\Delta V$  is  $< 5$  mm/s.  $\rightarrow$  Mutual escape is unlikely.

## Conclusions

- In our cratering model, ejecta imparts 6-12 times the momentum of the impactor.
- With this class of cratering event, kinetic impact is a viable method for the vast majority of threats.
- Even without cratering, kinetic impact is viable for large portion of threats, including keyhole deflection.
- We should start doing fast cheap scaled down kinetic impact missions to build our ability to predict and model cratering effects for various asteroid types.
  - Make it not so “uncontrolled”
- We should not shut down Arecibo. Radar gives:
  - Characterization which will allow better mission planning, results prediction
  - Precise measurement of deflection results
  - Longer lead times